

## IMPROVING THE MAXIMUM TRANSMISSION DISTANCE IN CV-QKD USING A NOISELESS LINEAR AMPLIFIER

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#### Can we increase this maximum distance ?



# OUTLINE

# I.Continuous-variable & coherent states QKD

# II. Heralded Noiseless Linear Amplifier (NLA)

# III.Improvement of CV-QKD performances with the NLA



# I.Continuous-variable & coherent states QKD



n

### **Discrete variables**

• Decomposition on a discrete basis  $|\psi\rangle = \sum c_n |n\rangle$ 

### Continuous variables

• Decomposition on a continuous basis  $|\psi\rangle = \int dx \ \psi(x) |x\rangle$ 

• Quadrature operators  $\hat{X}$  and  $\hat{P}$  = projection of the field's amplitude in the phase space, similar to the position and momentum for a massive particle

$$\hat{\boldsymbol{X}} = \left( \hat{\boldsymbol{a}} + \hat{\boldsymbol{a}}^{\dagger} 
ight) \sqrt{N_0}$$
  
 $\hat{\boldsymbol{P}} = \left( \hat{\boldsymbol{a}}^{\dagger} - \hat{\boldsymbol{a}} 
ight) i \sqrt{N_0}$ 



### DESCRIPTION OF A QUANTUM STATE

### Wigner function

Quasiprobability distribution



# Quadrature measurement

Homodyne detection





# GG02 PROTOCOL

PRL **88**, 057902 (2002) Nature **421**, 238 (2003)

### Quantum part

• Alice randomly selects  $x_A$  and  $p_A$  from a Gaussian distribution of variance  $V_A$ 

- The state  $|x_A+ip_A\rangle$  is sent to Bob
- ${\scriptstyle \bullet}$  Bob randomly measures the X or P quadrature





# GG02 PROTOCOL

### Equivalent Entanglement-Based version



Source of coherent states with: - amplitude proportional to  $\lambda(x_A + ip_A)$ - variance modulation  $V_A = \frac{1+\lambda^2}{1-\lambda^2} - 1$ R. Blandino - QCRYPT 2012

Quantum Info. Comput. **3**, 535–552 (2003)

RMP 84, 621 (2012)



LIMITS



### Can we increase this maximum distance ?





#### Maybe with an amplifier in Bob's station ?





#### Must add extra noise

Phys. Rev. D 26, 1817 (1982)



### **DETERMINISTIC PHASE SENSITIVE AMPLIFIER**



### Amplification of X

- ▶ Doesn't add extra noise → preserves the SNR
- Still amplifies the initial noise
- Only compensates homodyne imperfections

Journal of Physics B **42**, 114014 (2009)



# What happens if the amplifier is allowed to be non deterministic?



# II. Heralded Noiseless Linear Amplifier (NLA)



- $\bullet$  Performs the transformation  $|\alpha\rangle \rightarrow |g\alpha\rangle$
- Phase insensitive, but doesn't add extra noise
- Doesn't amplify the input noise
   R. Blandino QCRYPT 2012

T.C.Ralph and A.P.Lund, arXiv:0809.0326 (2008)



## Description of the NLA

- ► Cannot be unitary ➡ must be probabilistic
- Described by an unbounded operator  $g^{\hat{n}}$   $(g^{\hat{n}}|n\rangle = g^{n}|n\rangle)$

Transformation of usual states 
Gaussian operation

Coherent state

EPR state

 $g^{\hat{n}}|\alpha\rangle = e^{-\frac{|\alpha|^2}{2}} \sum_{n=0}^{\infty} \frac{\alpha^n}{\sqrt{n!}} g^n |n\rangle \propto |g\alpha\rangle$ 

 $g^{\hat{n}}|\lambda\rangle = \sqrt{1-\lambda^2} \sum \lambda^n g^n |n,n\rangle \propto |g\lambda\rangle$ n=0

Thermal state

 $g^{\hat{n}} \hat{\boldsymbol{\rho}}_{\rm th}(\lambda) g^{\hat{n}} = (1 - \lambda^2) \sum_{n=0}^{\infty} g^{2n} \lambda^{2n} |n\rangle \langle n| \propto \hat{\boldsymbol{\rho}}_{\rm th}(g\lambda)$ 



### Experimental implementation (quantum scissors)



- Output state truncated at 1 photon
- Good approximation for small amplitude



### Experimental implementation (quantum scissors)



► Good proof of principle R. Blandino - QCRYPT 2012 PRL 104, 123603 (2010)



### Quantum scissors



arXiv:0809.0326 (2008)

### **Polynomial approximation**

$$g^{\hat{n}} = \lim_{N \to \infty} \sum_{k=0}^{N} \frac{(\ln g)^k}{k!} \hat{\boldsymbol{n}}^k$$

PRA 80, 053822 (2009)



# VIRTUAL IMPLEMENTATION USING POST-SELECTION

#### Homodyne detection for Bob



N.Walk et al., arXiv:1206.0936 (2012)

#### Heterodyne detection for Bob



J. Fiurasek and N. J. Cerf, arXiv: 1205.6933 (2012)



#### **Theoretical implementation**

**Quantum scissors** T.C.Ralph and A.P.Lund, arXiv:0809.0326 (2008)

**Photon addition and subtraction** PRA 80, 053822 (2009)

**Experimental** implementation

#### **Quantum scissors**

PRL 104, 123603 (2010) Nat. Photon. 4, 316 (2010)

Photon addition and subtraction

Nature Photon. 5, 52 (2011)

Weak values arXiv:0903.4181 (2009)

Phase amplification PRA 81, 022302 (2010)

Short review of the experiments Laser Physics Letters 8, 411–417 (2011)

> Phase amplification Nat. Phys. 6, 767 (2010)



### Applications

**This talk** PRA 86,012327 (2012) **Error correction** PRA 84, 022339 (2011)

#### **Virtual implementation and QKD**

arXiv:1205.6933 (2012) arXiv:1206.0936 (2012)

#### **Optical loss suppression**

arXiv:1206.2852 (2012)

Cloning of coherent states PRA 86,010305 (2012)

**No violation of causality** PRA 86, 012324 (2012)



# III. Improvement of CV-QKD performances with the NLA



# 

# How does Bob's NLA improve

- the maximum distance of transmission?
- the maximum tolerable noise?
- the key rate?

# Assumption:

- Linear lossy and noisy Gaussian channel
  - transmittance T
  - added noise  $\epsilon$





# EFFECTIVE PARAMETERS

### Effective EPR parameter $\zeta$

$$\zeta = \lambda \sqrt{\frac{\left(g^2 - 1\right)\left(\epsilon - 2\right)T - 2}{\left(g^2 - 1\right)\epsilon T - 2}}$$

The NLA increases the entanglement •  $\zeta$  depends linearly on  $\lambda$ 

# Physical constraint

$$0 \le \zeta < 1 \Rightarrow 0 \le \lambda < \left(\sqrt{\frac{(g^2 - 1)(\epsilon - 2)T - 2}{(g^2 - 1)\epsilon T - 2}}\right)^{-1}$$



Effective transmittance  $\eta$  and noise  $\epsilon^g$ 

$$\eta = \frac{g^2 T}{\left(g^2 - 1\right) T \left[\frac{1}{4} \left(g^2 - 1\right) \left(\epsilon - 2\right) \epsilon T - \epsilon + 1\right] + 1}$$
$$\epsilon^g = \epsilon + \frac{1}{2} \left(g^2 - 1\right) \left(2 - \epsilon\right) \epsilon T$$

The NLA increases the transmittance and the noise
No dependence on  $\lambda$ 

# Physical constraints

$$\begin{cases} 0 \le \eta \le 1 \\ 0 \le \epsilon^g \end{cases} \} \Rightarrow \begin{cases} \epsilon \le 2 \\ g \le g_{max}(T, \epsilon) \end{cases}$$



# SECRET KEY RATE

### Key rate without the NLA

 $\Delta I(\lambda, T, \epsilon, \beta) = \beta I_{AB}(\lambda, T, \epsilon) - \chi_{BE}(\lambda, T, \epsilon)$ 

### Key rate with the NLA

= Key rate with the effective parameters without the NLA, weighted by the probability of success  $P_{\rm suc}$ 

$$\Delta I_{\rm NLA}(\lambda, T, \epsilon, \beta) = P_{\rm suc} \Delta I(\zeta, \eta, \epsilon^g, \beta)$$

Alice optimizes the variance modulation to maximize the key rate



# PROBABILITY OF SUCCESS

### **Optimistic** upper bound

$$P_{\rm suc} \le \frac{1}{g^2}$$

### Remarks on the probability of success

- Depends on the physical implementation
- Realistic probability of success may be much smaller
- Acts simply as a scaling factor
- Doesn't change the positivity or negativity of a key rate



- Strong losses regime  $(T << 1, \epsilon \neq 0)$ 
  - Without the NLA: minimum value of transmittance  $T_{\rm lim}$
  - With the NLA:

$$T_{\rm lim}^{\rm NLA} = \frac{1}{g^2} T_{\rm lim}$$

Tolerable losses are increased by  $20 \log_{10} g \, dB$ The maximum distance of transmission can be arbitrarily increased by increasing the gain



### IMPROVEMENT OF PERFORMANCES

### Maximized key rate





### IMPROVEMENT OF PERFORMANCES

### Maximum tolerable noise





### IMPROVEMENT OF PERFORMANCES

Can we arbitrarily increase the key rate? No...



- Optimal value of the gain
- If the gain if too important, the effective noise becomes too high
   R. Blandino - QCRYPT 2012



# CONCLUSION

# NLA in CV-QKD with a Gaussian lossy noisy channel

- Equivalent to an effective system without the NLA
- The maximum distance of transmission can be arbitrarily increased
- Improvement of the maximum tolerable noise
- Explicit formulas for GG02, same resultats for other CV-QKD protocols (same effective parameters)

### **Reference:** R. Blandino et al., PRA **86**, 012327 (2012)



# THANK YOU



# OUR TEAM









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